A DFT study of structural, electronic, and optical properties of Lead-free and Ge-based cubic perovskite RbGeX₃ (X= I, Br and Cl)

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The current research uses density functional theory (DFT) approximations in conjunction with the plane wave-pseudopotential method to investigate structural, electronic, and optical properties of Pb-free cubic perovskite RbGeX₃ (X= I, Br and Cl) materials. More specifically, Norm-conserving pseudopotential has been employed to describe the ion and valence electrons interaction, and GGA-PBE flavor is used to represent the exchange-correlation part of the energy of the GGA approximation. Lattice parameters obtained are 5.95, 5.55, and 5.29 Å for RbGeX₃ (where X=I, Br, and Cl), respectively, and they are comparable with available empirical and other values. The direct band gap nature of the three compounds under research here is shown and our values of the band gap energy \( E_g \) are agree with the other available results. Materials under research show responses to electromagnetic radiation starting from the infrared region to very high energies ( ~ 33 eV). The RbGeI₃ has the lowest \( E_g \) value at the low region energies and the highest optical response peaks but RbGeCl₃ has the highest optical response peaks at energies located near ~20 eV. Our results show that these materials are good candidates for photo electronic applications including solar cells.

1. Introduction

Perovskite materials are intensively under research and they are promising materials for a variety of photo optic applications because they have a tunable band gap due to the various elements which can be used to fabricate these materials[1-4]. These materials have been used to fabricate light-emitting diodes (LED) and photodetectors[5-8]. In the last decade, researchers have shown the potential of the organic-inorganic and all inorganic Pb-based perovskite materials for the next generation, low-cost solar cell. The efficiency of perovskite-based solar cells has increased dramatically up to 25.2% in 12 years[9]. The general formula of the halide perovskite is described as ABX₃, where A, B, and X are an organic or inorganic cation, a metal cation, and a halogen anion, respectively[10-13]. Lead replaced B in perovskites which contain the Pb element[14-18]. An issue here remains because of the toxicity of the Pb element. Researchers are trying intensively to fabricate and study an efficient solar cell from Pb-free perovskite material[19-28]. The toxicity issue can be dealt with by removing Pb and substituting an element such as tin (Sn) and germanium (Ge)[29-33].

Among Pb-free perovskites, Ge-based perovskites have significance in different applications including solar cells. Electronic properties are calculated by synthesizing AGeI₃ (where A= CH₃NH₃, HC(NH₂)₂, CH₃C(NH₂)₂, C(NH₃)₃, (CH₃)₂NH and (CH₃)₂C(H)NH₃)[34], it is shown that they have direct band gap and in the visible range spectrum. A highly luminescent (Up to 71%) Ge–Pb perovskite film with photoluminescence quantum efficiencies has been studied, and a new method is proposed for the eco-friendly light emitted technique[35]. A computational calculation combined with an experiment has confirmed that AGeI₃ (A=Cs, CH₃NH₃ or HC(NH₂)₂) has strong potential for future solar cells[36]. However, Ge-based solar cell has low power conversion and it has to be improved. The power efficiency is improved up to 0.57% for MAGeI₂Br₀.₃.

Organic cations as long as Cesium (Cs) can be replaced by Rb due to their similarity in electronic properties and their location in the periodic table. Recently researchers have investigated RbGeX₃ (where X= I, Br, and Cl) computationally and experimentally to provide more information on Ge-based perovskites. Density functional calculations have been performed to investigate structural, elastic, electronic, and optical properties of inorganic Ge-based perovskite with rhombohedral structure (space group R₃m) and cubic structure with space group Fm-3m. Results confirmed these compounds have a wide range of absorption in the visible and ultraviolet energy values which
confirms their future role in various optoelectronic applications and they are strong candidates for the future solar cell[32, 33]. Specially RbGeI₃ has the ability for harvesting maximum photon energies in IR and UV spectra[36].

Due to our perspective, more investigation is demanded to confirm findings on RbGeX₃ (where X= I, Br, and Cl). In the current work, structural, and optoelectronic properties are studied for the Ge-based RbGeX₃ (where X=I, Br, and Cl) in cubic structure (space group 221) using PBE-GGA and density fictional perturbation theory (DFPT) approximation.

2. Computational Methodology

In this research, structural properties, electronic properties, and optic response to electromagnetic radiation have been calculated in the framework of density functional theory (DFT)[39, 40]. For optic properties, Eigen values and Eigen functions resulted from density functional perturbation theory (DFPT) calculations used in Kramer-Kronig relations[41]. All calculations have been performed with approximations implemented in ABINIT software[42].

Pseudopotential as core-valence states interaction and plane-wave as basis sets are chosen to perform calculations. Pseudopotentials used here are previously generated by M.Krack where Goedecker, Teter, and Hutter (GTH) approach was used[43]. The electronic configuration of each atom here is as follows: Rb: 4s⁵ 5p⁶, Ge: 4s⁴ 4p² I: 5s² 5p⁵, Br: 4s² 4p⁶, Cl: 3s² 3p⁵. In the beginning, convergence calculations of the total energy (Etolerance = 2.7 × 10⁻¹¹ eV ) of the compounds have been done on Monkhorst–Pack mesh[44] of k-points and cut-off energy (Ecut-off). The Monkhorst–Pack special k-points were 6x6x6 for structural and electronic properties and set to 24x24x24 to calculate Eigen values and Eigen functions needed for optic properties calculations. The Etot converged when Ecut-off was 1235 eV. Full optimization for the structure dimensions has been performed for different pressures using PBE-GGA approach and (BFGS)[45, 46].

The lattice dimensions parameters and Wyckoff positions were obtained and used for later calculations. Band structure energies for all compounds are calculated employing PBG-GGA. As is well known, the PBE-GGA method underestimated the value of the band gap energy Eₔ when compared to the experimental values due to the strong fermion interactions in the heavy physical systems[47]. By performing DFPT, Eigen values and Eigen functions are obtained and used to calculate both the imaginary part and real part of frequency-dependent conductivity σ(ω):

\[ σ(ω) = σ₁(ω) + iσ₂(ω) \]  

Where σ₁(ω) and iσ₂(ω) are real and imaginary parts of conductivity, respectively, and ω is frequency. The real part can be calculated using Kubo-Greenwood (KG) formulation[48]:

\[ σ₁(ω) = \frac{2\pi}{n} \sum_i |F_{ij}|^2 \delta(\epsilon_i - \epsilon_j - \omega) \]  

Where Ω is atomic volume, Fij is Fermi-Dirac distributions and e is the dielectric function. And the imaginary part is calculated using Kramers-Kronig relations:

\[ σ₂(ω) = \frac{2}{\pi} \int P \frac{σ₁(\nu)ω}{(ω^2 - \nu^2)} dν \]  

Where P is the principal value of the integral. The real part of the frequency-dependent dielectric response ε₁(ω) and imaginary parts are calculated from imaginary and real parts of conductivity, respectively:

\[ ε₁(ω) = 1 - \frac{4\pi}{ω} σ₂(ω) \]  
\[ ε₂(ω) = \frac{4\pi}{ω} σ₁(ω) \]  

The total dielectric faction is the sum of the two parts and is related to the real part of the refractive index n(ω) and imaginary part k(ω) as follows:

\[ n(ω) = \frac{1}{2}\sqrt{|ε(ω)| + ε₁(ω)} \]  
\[ k(ω) = \frac{1}{2}\sqrt{|ε(ω)| - ε₁(ω)} \]  

Refractive index components are related to the reflectivity r(ω) and absorption coefficient α(ω) through the equations below:

\[ r(ω) = \frac{|1 - n(ω)|^2 + k(ω)^2}{|1 + n(ω)|^2 + k(ω)^2} \]  
\[ α(ω) = \frac{4\pi}{n(ω)} σ₁(ω) \]  

3. Results and Discussions

3.1 Structural properties

The present research considers the ideal cubic of the RbGeX₃ (where X=I, Br, and Cl) perovskite which consists of a polyhedron that shares 12 halide atoms on the corners and is located inside the cube and one Rb atom in the body and 8 Ge atoms at the corners of the cube (figure 1). The unit cell contains one molecule of five atoms of Wyckoff positions as Rb 1b (0.5, 0.5, 0.5), Ge 1a (0.0, 0.0, 0.0), and X 3d (0.0, 0.0, 0.5).

Figure 1: Ideal cubic structure of RbGeX₃ (X= I, Br, and Cl).

Lattice constant (a), the bulk modulus (B), and its derivative (B’), are calculated by fitting the change in energy (Etot) as a function of varied volumes (V) by the Murnaghan-Birch equation of state as shown below[49, 50]:
Here, the first Brillouin zone (IBZ) and high symmetry k-points path \( \Gamma-X-M-R-\Gamma \) are chosen to represent band structure energies (Figure 3). The Fermi level is set to zero. Band structure shows that the compounds RbGeX3 (where X=I, Br, and Cl) have the conduction band minimum (CBM) and valence band maximum (VBM) at R symmetry k-point (direct band gap value).

![Figure 3: Band structure of RbGeX3 (where X=I, Br, and Cl) compounds along high symmetry points path (Gamma-X-M-R-Gamma) using GGA approximation.]

It is widely known that the PBE-GGA approximation underestimates the values of the \( E_g \). The resulting band gap energy values \( E_g \) from GGA approximation of the three halide perovskites of this work and other works results are listed in table 2. It is noted that results here comparably agree with other results which are also calculated using the GGA method. The band gap values are 0.53, 0.61, and 0.94 eV for RbGeI3, RbGeBr3, and RbGeCl3, respectively. The \( E_g \) increases ongoing from I to Cl due to the weakening of the interaction of the electrons in the valence and conduction bands which depends on the number of electrons in both bands. These results show the tunable band gap characteristic of this group of perovskite materials as the halogen atom changes.

![Figure 2: Total energy in eV as a function of volume in Å³.]

**Table 1:** Calculated lattice constant values (a), Bulk modulus (B), and its derivative (B’) of cubic RbGeX3 (X= I, Br, and Cl).

<table>
<thead>
<tr>
<th>Materials</th>
<th>( a ) (Å)</th>
<th>Experimental</th>
<th>Other works</th>
<th>( B ) (GPa)</th>
<th>Other work</th>
<th>( B’ ) (GPa)</th>
<th>Other work</th>
</tr>
</thead>
<tbody>
<tr>
<td>RbGeI3</td>
<td>5.95</td>
<td>5.99</td>
<td>5.96 ( ^a ) 5.91 ( ^c )</td>
<td>17.76</td>
<td>12 ( ^b )</td>
<td>4.45</td>
<td></td>
</tr>
<tr>
<td>RbGeBr3</td>
<td>5.55</td>
<td></td>
<td>5.53 ( ^c ) 5.57 ( ^d )</td>
<td>22.74</td>
<td>22.48 ( ^d )</td>
<td>4.31</td>
<td>4.14 ( ^d )</td>
</tr>
<tr>
<td>RbGeCl3</td>
<td>5.29</td>
<td></td>
<td>5.27 ( ^c ) 5.31 ( ^d )</td>
<td>26.71</td>
<td>26.15 ( ^d )</td>
<td>4.016</td>
<td>4.05 ( ^d )</td>
</tr>
</tbody>
</table>

**Table 2:** Calculated band gap values \( E_g \) in eV RbGeX3 (X= I, Br, and Cl) calculated in GGA approximation.

<table>
<thead>
<tr>
<th>Materials</th>
<th>( E_g ) (eV)</th>
<th>This work</th>
<th>Other works</th>
</tr>
</thead>
<tbody>
<tr>
<td>RbGeI3</td>
<td>0.53</td>
<td></td>
<td></td>
</tr>
<tr>
<td>RbGeBr3</td>
<td>0.61</td>
<td></td>
<td>0.67 ( ^a )</td>
</tr>
<tr>
<td>RbGeCl3</td>
<td>0.94</td>
<td></td>
<td>0.96 ( ^a )</td>
</tr>
</tbody>
</table>

\( ^a \) reference \( ^{[37]} \),
To explain the CBM and VBM bands creation, the total and partial density of the state of RbGeX₃ (where X=I, Br, and Cl) compounds are calculated using GGA approximation and plotted as it is shown in figure 4. The valence state is created due to the contribution of the s states of the Ge and p states of the halogen atom and Ge atom. While the conduction band was created as the result of the contribution of the p states of the Ge and halogen (X) atoms. The hybridization of the p states of Ge and halogen atoms indicates the covalent nature of the bond between them.

The relationship between Im(\(\epsilon(\omega)\)) and \(E_g\) value is obvious. It is shown in Figure 5a that the response of the three compounds RbGeX₃ (where X=I, Br, and Cl) is similar. The first transition (region I, fig.5a) happens from VBM to the CBM with represents the critical point and corresponds the \(E_g\) value. The response behavior can be explained in conjunction with DOS. The response of the RbGeI₃ is the highest because of the p state of the halogen contribution to the creation of the VBM and the iodine has the highest number of electrons. Region II is due to the transition of electrons between s state of Ge and d state of the Rb and the region III response appears as the result of transitions of the electron from s state of halogens and d state of the Rb element. It is also worth mentioning that the response shifts toward lower energies as the band gap decreases and the \(E_g\) value has an inverse relationship with Im(\(\epsilon(\omega)\)).

One of the other significant optical properties is the real part of the dielectric function Re(\(\epsilon(\omega)\)). Figure 5b shows the response of the Re(\(\epsilon(\omega)\)) to electromagnetic radiation ranging from 0 to 34 eV. The Re(\(\epsilon(\omega)\)) response of the three compounds RbGeX₃ (where X=I, Br, and Cl) is very similar in pattern but different in values. The response starts to appear at zero frequency Re(\(\epsilon(\omega)\)) at a certain point which represents the dielectric constant and increases to reach a maximum value at energies between 2 to 3 eV then decreases until it goes below 1 at energies ~7.5 to 9.3 eV. The dielectric constants of the three compounds are 7.21, 5.45, and 4.32 for RbGeI, RbGeBr, and RbGeCl, respectively.

The real part of the refractive index is also plotted here (Figure 5c), the behavior of the refractive index of the three compounds is similar to the Re(\(\epsilon(\omega)\)). Except for the point where the Re(\(\eta(\omega)\)) starts to respond represents the actual refractive index value of the compound. The refractive index of the three compounds is 2.67, 2.33, and 2.07 for RbGeI₃, RbGeBr₃, and RbGeCl₃, respectively. The refractive index increases until reaches its maximum value at energies ranging between 3.32 to 2.38 eV then decreases until it goes below the unity of the three compounds at energies ranging from 9.63 to 14.6 eV because the phase velocity of the light becomes greater than the speed of light itself and these materials show the superluminality in this, particularly range of energy[52].

The absorption coefficient \(\alpha\) (m⁻¹) starts to increase at energies in the infrared region until it reaches its maximum value at energies ranging from 18.9 to 20.04 eV for RbGeX₃ (where X=I, Br, and Cl) ongoing from I to Cl, RbGeCl₃ has the maximum value of \(\alpha\) (m⁻¹) (2.07×10⁵ m⁻¹) followed by RbGeCl₃ (1.87×10⁵ m⁻¹) and RbGeI₃ (1.67×10⁵ m⁻¹) (figure 5d).

The imaginary part of the dielectric function Im(\(\epsilon(\omega)\)), the Real part of the dielectric function (Re(\(\epsilon(\omega)\))), Im(\(\epsilon(\omega)\)) and Re(\(\epsilon(\omega)\)) parts and absorption coefficient \(\alpha\) (m⁻¹) of cubic RbGeX₃ (X=I, Br and Cl) perovskite are carried out here using a

**Conclusion**

Structural, opto-electronic properties of ideal cubic and Pb-free RbGeX₃ (X=I, Br, and Cl) perovskite are carried out here using a
plane-wave-pseudopotential approximation in conjunction with the DFT framework. Lattice parameters obtained here are 5.95, 5.55, and 5.29 Å for RbGeX₃ (where X=I, Br, and Cl). Our values of α agree well with the available experimental values and other values except for the bulk modulus of RbGeI₃ where it is ~38% larger, the benefit to our result due to the reasonable variance trend of our results. The lattice parameter values show a decrease with increasing ionic radius of the halide dopant in the cell (i.e., from I to Cl). The band gap values are 0.53, 0.61, and 0.94 eV for RbGeI₃, RbGeBr₃, and RbGeCl₃, respectively. The $E_g$ increases from I to Cl due to the weakening of the interaction of the electrons in the valence and conduction bands, which depends on the number of electrons in both bands. These materials start to respond to the electromagnetic spectrum in the IR region and end at energy near 34 eV. RbGeI₃ is a better absorber and retainer to the radiation in low energies and RbGeCl₃ is a better absorber at medium and high energies. At the highest point where these materials can respond, RbGeI₃ has a better response again. Due to the lack of theoretical studies of these materials, and the absence, as to our knowledge, of experimental studies, more investigations needed to be done on these materials.

Conflict of interests
None

Contribution
The author certifies that he has worked sufficiently in the work to take public responsibility for the content, including working on the concept, design, analysis, writing, or revision of the manuscript. Furthermore, the author certifies that this material or similar material has not been and will not be submitted to or published in any other publication before its appearance in the Passer journal of basic and applied science.

References


